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Introduction

Sommerfeld Integrals (SIs) are a type of oscillatory integral that is used to describe the spatial Green’s functions of dipole sources radiating in layered media. They may be used in conjunction with moment methods to model scattering and radiation in such media.

The Complex Image (CI) method is a rapid but approximate technique used to calculate spatial Green’s functions without recourse to time-consuming numerical integration of SIs, thus greatly decreasing the fill time for the moment method interaction matrix.

The CI method operates by approximating part of the spectral Green’s function as a sum of complex exponentials, thereby allowing one to inverse transform it in to a closed-form spatial Green’s function. This process is generally an order of magnitude faster than numerical integration.

The main problem with the CI method is that there is no way of checking the accuracy of the spatial Green’s functions it is used to calculate, until they have been derived and compared with established techniques—like numerical integration or asymptotic evaluation. This paper introduces for the first time a new variant of the CI method that uses a new path in the complex plane.

2 Modeling a VMD above a Half-space

Consider the case of an infinitesimal VMD in free space radiating above a half-space of permittivity $\varepsilon = \varepsilon_0$. The vector potential dyadic Green’s function for this problem geometry is $G_{\mu} = i\varepsilon_0 G$, where:

$$G_{\mu}(p, z; p', z') = \frac{1}{4\pi} \int_{\mathbb{R}^3} \hat{G}_{\mu}(p, z; \hat{z}')(k_p | p - \hat{z}|) dpdz,$$

(1) where $G_{\mu}$ is the Green’s function and $k_p$ is the wave-number in the $p$ direction.

The spectral Green’s function $\hat{G}_{\mu}$ is:

$$\hat{G}_{\mu}(k_p, z; \hat{z}') = \frac{1}{\sqrt{2\pi}k_p} \left[ e^{-ik_p|z-\hat{z}'|} + R^{TE} e^{-ik_p(|z-\hat{z}'|+\theta_0)} \right],$$

(2) where the wave-number in free space $k_0 = \frac{\omega}{c}$.

The TE plane-wave spectral reflection coefficient $R^{TE}$ is:

$$R^{TE} = \frac{\varepsilon_{||} - k_0^2}{\varepsilon_{||} + k_0^2},$$

(3)
2.1 Asymptotic Model

As \( t \) and \((t+s')\) increase asymptotic techniques can be used to derive expressions for \( G^{+}_m \) such as that given in equation 4:

\[
G^{+}_m = -\frac{e^{-j\mu s'}}{4\pi\mu} - \frac{e^{-j\mu s'}}{4\pi\mu} \frac{\partial G^{+}_m}{\partial s'} \bigg|_{s' = 0} + G^{+}_m, \tag{4}
\]

where

\[
\alpha_1 = \sqrt{\frac{\mu}{\epsilon}} \frac{\mu \mu_0}{\epsilon_0} (1 + \epsilon)^2
\]

and

\[
h_m = h_0 \frac{\mu_0}{\sqrt{\epsilon_0 + \epsilon(1 + \epsilon)^2}}. \tag{6}
\]

The first term in equation 4 is the direct term of the VMD and is obtained by analytical integration.

The second term is the reflected term derived using the method of stationary phase [1], a single asymptotic technique that is valid when the source and observer are not close to the interface.

These two terms represent two optical waves—one emanating from the source and one from its image (which is modified by \( R_{in}(y_{imag}) \)).

From equations 3 and 6 it can be seen that as \( (t+s') \to \infty \) and \( \mu \to 0 \), \( R_{in}(y_{imag}) \) tends to \(-1\).

Therefore the first two terms in equation 4 cancel each other out, resulting in the field along this interface region being dominated by \( G^{+}_m \).

\( G^{+}_m \) may be derived using complicated and problem dependent asymptotic techniques and represents a lateral wave [1]. Lateral waves are inhomogeneous waves that have an algebraic decay of \( jn \) in the \( y \) direction, and are evanescent at some nomes away from the interface.

2.2 Standard Complex Image Model

The standard CI method proceeds by parametrically sampling the TE plane-wave spectral reflection coefficient \( R_{in} \) along Path 1 (see fig. 1) in the \( k_{ow} \) plane [2]. The parameter \( T \) is chosen so that Path 1 avoids the branch points at \( k_0 = k_x \) and \( k_0 = k_x \) marked with an x in fig. 1.)

![Image of Complex k Plane](image)

Figure 1: Sample Paths on the Complex k Plane.
The samples of $\tilde{a}_n^k$ are then approximated by a sum of $N$ complex exponentials of a real variable, using a system identification technique such as the SVD-Rpty method [3, 4] (any error in this approximation introduces an error into the real Greens' Function, which is then transformed back into a function of the complex variable $a_n^k$ [1]).

$$\tilde{a}_n^k = \sum_{m=1}^{N} a_n^m e^{-i \omega_b^m}. \quad (7)$$

An approximate closed-form expression for $G_z^R$ can now be derived using equations 1, 7 and the Sommerfeld identity [11]:

$$G_z^R = \frac{e^{i \omega_b^m b_0}}{4\pi b_0} - \frac{e^{-i \omega_b^m b_0}}{4\pi b_0} \sum_{m=1}^{N} a_n^m e^{-i \omega_b^m b_0} \quad (8)$$

where

$$R_n = \sqrt{\mu - a^2} + \sqrt{\mu + a^2} \quad (9)$$

and

$$R_n = \sqrt{(p - j)^2 + (\mu + a^2)} \quad (10)$$

When compared with numerical integration, the standard CI method only produces accurate results for a limited range of source and observer coordinates; for example, in Fig. 2, it fails at $\mu = 0$, i.e. for $\mu < 0.25$. Various complex image terms were used to approximate $\tilde{a}_n^k$.

The inaccuracy of the standard CI method in this example can be explained as follows: the terms in equation 8 represent a spherical wave emanating from the source and a summation of spherical waves emanating from complex locations in the lower half-space, modified by complex residues. In the spatial domain, the CI method therefore characterises the source's interaction with the interface as a summation of spherical waves; these spherical waves are not, however, an ideal basis to model a lateral wave.

The standard CI method works well enough in regions where the lateral wave is not strongly excited, such as when the source and/or observer are far from the interface. However, it fails when one moves along the interface in the $\mu$ direction, because $\mu$ gradually dominates the poles in equation 10.

$$\mu \rightarrow \infty \Rightarrow R_n \rightarrow 0. \quad (11)$$

The information about $R_n^k$ that is modelled by the poles in equation 7 is thus lost, causing significant errors in the spatial Greens' function.

2.3 New Complex Image Model

The new CI method proposed in this paper samples $\tilde{a}_n^k$ along Path 2 in the complex $\lambda_{\mu}$ plane (see Fig. 1). Path 2 starts at the branch point at $\lambda_{\mu} = \lambda_{\mu}^*$ and terminates just off the imaginary axis near the branch point at $\lambda_{\mu} = \lambda_{\mu}$. This path concentrates the approximation process on the behavior of $\tilde{a}_n^k$ along the portion of the $\lambda_{\mu}$ plane that describes waves that excite lateral wave modes. Therefore the new technique is much more robust along the interface (see Fig. 2).

This new CI method shows that the integration of equation 1 can be performed over a limited range of the variable $\lambda_{\mu}$ to improve the accuracy and robustness, in a predetermined region of space, of the CI spatial Greens' function $G_z^R$. 

848
3 Conclusions

The far-field asymptotic expressions for the vector potential Greens' function of an infinitesimal Vertical Magnetic Dipole (VMD), radiating above a half-space, are used to explain the failure of the standard CI method, when the source and observer are near to the interface—but separated by a large lateral distance. In these types of problems the field at the observation point is dominated by a localized wave that is not modeled accurately by the standard CI method. It is shown for the first time that by applying the CI approximation process over a different path in the k_0 plane, the accuracy and robustness of the CI method for such a problem geometry can be improved.

References


