Exploring Autoionization and Photo-Induced Proton-Coupled Electron Transfer Pathways of Phenol in Aqueous Solution

Thomas A. A. Oliver, Yuyuan Zhang, Anirban Roy, Michael N. R. Ashfold, and Stephen E. Bradforth

1 School of Chemistry, University of Bristol, Bristol, BS8 1TS, UK, 2 University of Southern California, CA 90089 USA

AUTHOR INFORMATION

Corresponding Author

* email: stephen.bradforth@usc.edu. Tel: (213) 740-0461. Fax: (213) 740-3972

† Department of Chemistry and Biochemistry, Montana State University, Bozeman, MT 59717, USA
‡ These authors contributed equally to this work
ABSTRACT The excited state dynamics of phenol in water have been investigated using transient absorption spectroscopy. Solvated electrons and vibrationally cold phenoxyl radicals are observed upon 200 nm and 267 nm excitation, but with formation timescales that differ by more than four orders of magnitude. The impact of these findings is assessed in terms of the relative importance of autoionization versus proton-coupled electron transfer mechanisms in this computationally tractable model system.

KEYWORDS UV photochemistry, solvated electrons, proton transfer, electron transfer, non-radiative decay
Phenol is the chromophore of the amino-acid tyrosine, a key residue in ground and excited state biochemical redox reactions, such as the reaction center of Photosystem II (PSII), and enzymatic catalysis. The photophysical properties of the phenol chromophore in aqueous solution have been extensively studied on both the ground and excited potential energy surfaces (PESs). Previous flash photolysis and transient absorption studies of phenol and tyrosine in aqueous solution reveal the formation of phenoxyl (or tyrosyl) radicals and solvated electrons, but the precise details of the mechanism(s) leading to these photoproducts have remained elusive. The debate as to whether the process involves excited state proton transfer (ESPT), autoionization, O–H bond fission (H-atom transfer), or excited-state proton-coupled electron transfer (PCET) pathways remains a matter of some controversy. A PCET pathway is understood to involve the concerted transfer of a proton and electron, rather than a stepwise transfer, and is recognized as an important mechanism in the functioning of many biological systems. To date, electronic structure calculations have only examined small phenol–water clusters, revealing that direct photoexcitation to the \( ^1\pi\sigma^* \) state leads to a net transfer of an H atom to the water cluster.

Here we report the use of ultrafast transient absorption (TA) spectroscopy methods to investigate the excited-state dynamics of phenol-\( h_6 \) (in H\(_2\)O) and its phenol-\( d_1 \) isotopologue (in D\(_2\)O) using a femtosecond UV pump excitation, and a broadband white light super-continuum probe that spans the electronic absorption of the phenoxyl radical and solvated electron species. The experimental apparatus has been described previously, and is summarized in the supporting information (SI).

Fig. 1(a) displays TA spectra for phenol-\( h_6 \) in H\(_2\)O at the displayed time delays, \( t \), following 200 nm (6.2 eV) excitation. Two notable features are evident after \( t = 200 \) fs, the time-resolution
Figure 1. (a) TA spectra of 18 mM phenol-$h_0$/H$_2$O solution for displayed pump-probe time delays and 200 nm excitation (b) Early time ($t \leq 1$ ps) TA spectra displayed on an expanded scale.

of our transient absorption experiments at this excitation wavelength: vibrationally cold phenoxy radicals with an electronic origin at $\sim400$ nm, and a broad absorption attributed to solvated electrons (that peaks at $\sim700$ nm). The profile of the broad solvated electron transient evolves and its band centre shifts to shorter wavelength during the first 2 ps, as the electrons become fully hydrated. The assignment of the solvated electron transient was confirmed by the addition of HCl, a known electron scavenger (see SI).

Expanded views of early time ($t \leq 1$ ps) TA spectra in the 350–500 nm probe region are shown in Fig. 1(b). These data reveal an additional feature centered at 425 nm (indicated by the arrow), which we assign based on prior literature, to the ground state phenol radical cation (PhOH$^+$), and which disappears on a sub-picosecond timescale. To the best of our knowledge, this represents the first observation of the PhOH$^+$ species in aqueous solution near neutral pH.

The vertical ionization potential (VIP) of phenol (PhOH) in water is 7.9 eV relative to the vacuum level. Excitation at 200 nm (6.2 eV) is unable to induce vertical ionization of phenol. The thermodynamic threshold for producing fully solvated charged products, i.e., the phenoxy radical (PhO$^-$), a proton and a solvated electron, is estimated to be in the range 4.3 – 4.5 eV (see SI for the free energy landscape of PhOH and associated photoproducts). The observation of the
PhOH$^+$ reaction intermediate (Fig. 1(b)) thus points in favor of an autoionization mechanism, followed by rapid deprotonation. The quantum yield of solvated electron formation was extracted via calibration using the known solvated electron yield for photodetaching an electron from hydroxide at 200 nm, yielding values of 26% (at $t = 1$ ps) and declining to ~14% at $t = 800$ ps. We attribute this decline to geminate recombination between PhO$^-$ and solvated electrons, and may involve the third geminate partner, the proton, as discussed in greater detail later. Using known molar extinction coefficients, the ratio of phenoxyl and solvated electron products was found to be very close to 1, suggesting few, if any, phenol molecules undergo homolytic O–H bond photodissociation at 200 nm.

These TA results are strikingly different to those reported previously for phenol-$h_{\nu}$ in the non-polar solvent, cyclohexane, wherein 200 nm excitation was seen to yield prompt vibrationally hot phenoxyl radicals (within 200 fs), with no evidence for solvated electrons. Such a finding is consistent with direct photodissociation on a repulsive $^1\pi\sigma^*$ PES, as found in gas phase experiments. The solvent dependent photochemistry of phenol, is likely dictated by the magnitude of the hydration contribution to the thermodynamics of the charged reaction species. Water is energetically able to support charged products such as protons and solvated electrons, whereas non-polar solvents like cyclohexane are unable to do so, therefore phenoxyl radical generation in non-polar solvents largely mirrors the gas phase reaction.

Our recent transient absorption studies of para-methylthiophenol ($p$-MePhSH) in solution illustrate that the electrostatic properties of the solvent are certainly not the sole factor dictating the photoinduced dynamics in solution. These studies showed no evidence for autoionization or PCET pathways when exciting at 267 nm or at 200 nm, even in polar ethanol solution which can support and stabilize both the electron and dipolar $p$-MePhS radical; 'gas-
phase’-like (i.e. S–H bond fission) dynamics were seen to prevail. Two conclusions follow: First, the fact that p-MePhSH in ethanol has a lower VIP (7.2 eV) than PhOH in water (7.9 eV) shows that a lower VIP does not necessarily predicate a larger autoionization yield. Second, we conclude that the O–H bond fission process as seen for PhOH in aqueous solution must be far slower, and thus a less competitive non-radiative decay pathway, than the rate of autoionization when excited at 200 nm.

We now focus on the dynamics of phenol-\(h_6\) following excitation at 267 nm, which populates the lower \(S_1(1\pi\pi^*)\) state. Spectra recorded at \(t < 1\) ns are dominated by ESA from the \(S_1\) state and a growing feature peaking at \(~600\) nm that has been observed previously and assigned to ESA of phenol excimers. No spectral signatures attributable to phenoxy radicals or solvated electrons are discernible in TA spectra recorded at \(t < 1\) ns (see SI). The early time kinetics of the TA signal at \(~600\) nm (at the peak of the transient) are unaffected by addition of electron scavengers (see SI), implying that the quantum yield of electron ejection is negligible at \(t < 1\) ns. At \(t > 2\) ns, the absorption maximum shifts to longer wavelength, as the phenol excimer signal decays and the broader feature that we assign to the solvated electron rises (See Fig. 2(a)).

![Figure 2](image)

**Figure 2.** Comparison of (a) 90 mM phenol-\(h_6\)/H\(_2\)O and (b) phenol-\(d_1\)/D\(_2\)O TA data following 267 nm excitation, measured at selected time delays. Insets display the same time delays but focus on the 375-650 nm probe region.
The TA spectrum recorded at $t = 14$ ns clearly shows the signature of the phenoxy radical, and electron scavenging experiments (Fig. 3(a)) confirm the presence of solvated electrons, and the small remaining transient at $\sim 600$ nm is due to residual phenol excimers. The phenoxy radical yield was found to be 13% at $t = 13$ ns via calibration using the known PhO$^\cdot$ (plus solvated electron) photodetachment yield from phenolate.$^{22}$ From literature molar extinction coefficients,$^{22,24}$ we then established that for 267 nm excitation of phenol in aqueous solution, solvated electrons and phenoxy radicals were produced in a 1:1 ratio. The yield of phenoxy radicals increases from the earliest delay time to $\sim 14$ ns. We, however, cannot ascertain the time in which the yield maximizes due to the limited number of discrete pump-probe time delays measured.

**Figure 3.** (a) TA spectra of phenol with and without 50 mM HCl (the electron scavenger) measured at $t = 13$ ns. (b) TA spectra of phenol-$h_6$/H$_2$O and phenol-$d_1$/D$_2$O measured at $t = 13$ ns.

TA spectra recorded at selected pump-probe time delays following 267 nm excitation of phenol-$d_1$ in D$_2$O are shown in Fig. 2(b). These were acquired sequentially to the phenol-$h_6$ in H$_2$O data in order to minimize changes in experimental conditions. These data display the same spectral features and kinetics as for the case of phenol-$h_6$ in H$_2$O. Comparing the TA spectra from phenol-$h_6$/H$_2$O and phenol-$d_1$/D$_2$O measured at $t = 13$ ns (overlaid in Fig. 3(b)) also shows negligible difference in the relative yields of phenoxy radicals and solvated electrons in the two
cases, implying a negligible kinetic isotope effect \(i.e.\ KIE = 1.0 \pm 0.4\). Note, that within the upper and lower bounds of the error \((0.6–1.4)\), the returned KIE is very small, and does not change our discussion.

We can envision several possible mechanisms for forming phenoxyl radicals and solvated electrons on a nanosecond timescale following excitation at 267 nm; photodissociation of the O–H bond, as per the gas phase, excited state proton transfer followed by electron detachment, PCET or autoionization. In turn, we consider the merits of each potential mechanism.

Previous experiments in non-polar solvents \(^{33}\) and in the gas phase \(^{33,39,44,45}\) have shown that PhOH molecules excited to the \(S_1\) (\(1\pi\pi^*\)) state can tunnel beneath a conical intersection (CI) associated with the \(S_2/S_1\) (\(1\pi\sigma^*/1\pi\pi^*\)) diabatic states on a nanosecond timescale, resulting in homolytic O–H bond fission, vibrationally excited PhO\(^-\) radicals and translationally excited H atoms. For this to be a viable route to forming the solvated electrons observed in the present study, which appear on a very similar nanosecond timescale, would require rapid H atom fission into H\(^+\)(aq) and e\(^-\)(aq) after the bond scission. This seems very improbable, given that the equilibrium constant for this process favors neutral H atoms rather than ionic products (by 6 orders of magnitude); \(^{26,46}\) recall that reaction with H\(^+\) is the basis for the e\(^-\) scavenging process observed in Fig. 3(a). The corresponding gas phase studies involving phenol-\(d_6\) show O–D bond fission does not occur, \(^{35}\) because the tunneling probability for a D atom under the \(S_2/S_1\) CI is greatly reduced. \(^{17,39}\) Therefore, such an explanation is also unable to account for the measured KIE \((1.0 \pm 0.4)\). The absence of any homolytic O–H bond fission contribution likely reflects a third effect of the polar water on the reaction: the barrier to dissociation under the \(S_2/S_1\) CI must be larger (or significantly altered) in aqueous solution \(cf.\) in cyclohexane \(^{33}\) or in the gas phase \(^{39,47}\) – as predicted by previous calculations for phenol–H\(_2\)O clusters. \(^{16}\)
Many prior studies have explored ESPT in phenol.\textsuperscript{7,48-51} The pKa in the $S_1$ state is only 4 (cf. 9.8 in the $S_0$ state),\textsuperscript{52} and potentially enables proton transfer from the UV-excited phenol to proximal water molecules and formation of an excited-state phenolate anion that can eject an electron to solvent. If the second step is rate limiting, we would expect to observe the ESA band of the phenolate($S_1$) anion at ~500 nm with a lifetime of 22 ps concomitant with electron appearance on a similar timescale.\textsuperscript{22} Neither is consistent with the spectroscopy or kinetics evident in the present TA study. Alternatively, an ESPT process wherein the initial deprotonation step was rate limiting should be expected to have a large associated KIE. For example, naphthol derivatives have ESPT KIEs ranging between 1.7 and 3.8.\textsuperscript{53} Such behavior is not observed in our experiments either. Further support for excluding ESPT as the source of the observed charged photoproducts is provided by previous flash photolysis experiments which concluded that <3% of excited phenol molecules underwent ESPT within the $S_1$ lifetime,\textsuperscript{47} and several other reported failures to observe any signatures of ESPT for phenol in aqueous solution.\textsuperscript{48,49}

We now turn to consider possible PCET based explanations. PCET on the ground state potential (i.e. following internal conversion from $S_1$) is considered unlikely given the observed lack of any kinetic isotope effect (KIE = 1.0 ± 0.4), cf. the KIE = 2.5 ± 0.7 value derived from electrochemical studies of PCET involving phenol($S_0$) molecules by Costentin \textit{et al.}.\textsuperscript{4}

Several previous studies have demonstrated PCET on an excited state PES (albeit typically occurring on picosecond timescales \textsuperscript{54-56}). Photoinduced PCET thus merits discussion in the context of phenol. The electron and proton could be transferred into bulk solution on similar timescales in a concerted fashion, but the electron would also have to escape the first solvation shell in order to avoid fast inner-cage recombination with the phenoxy radical and be sufficiently separated to account for the deduced ~13% phenoxy radical yield at $t = 13$ ns.
Combined Monte Carlo and quantum chemical calculations show a PhOH—OH$_2$ hydrogen bond length of ~1.8 Å in the PhOH($S_1$) state. The electrons and protons would need to be ejected distances corresponding to several solvent shells (i.e. > ~3.6 Å) from the parent molecule, and from each other, in order to reduce loss by diffusive recombination sufficiently to achieve the observed radical and solvated electron survival probabilities. Studies by the Hammes-Schiffer group indicate that the KIE of ground state PCET reactions increases with the distance between acceptor and donor, and that the rate is thus determined by the solvent re-organization timescale rather than PCET. It is unclear whether such arguments will be similarly applicable to excited state PCET but, in the absence of supporting theory, we deem it unnecessarily complicated to suggest excited state PCET as the dominant route to charged products following 267 nm photoexcitation of phenol.

A simpler explanation assumes that the same mechanism identified following excitation at 200 nm applies when exciting at 267 nm also, but that it occurs on a much slower timescale. 267 nm excitation promotes phenol molecules to vibrationally excited levels (predominantly ring-breathing vibrations) of its $S_1$ state which relax to the vibrational ground state ($S_1(v=0)$) in several picoseconds – thereby reducing the total energy to ~4.5 eV. The PhO$^+$($aq$) + H$_3$O$^+$($aq$) + e$^-$($aq$) asymptote lies between 4.3–4.5 eV (see SI), i.e. only just below that of a PhOH($S_1(v=0)$) molecule. The deduced nanosecond reaction timescale could thus simply reflect slow autoionization from the $S_1$ origin to the PhOH$^+$($aq$) + e$^-$($aq$) asymptote, followed by rapid deprotonation. The non-observation of phenol radical cations can be rationalized on the basis of their very short lifetime, recall the sub-picosecond decay of the PhOH$^+$ transient we observe at 200 nm (see Fig. 1(b)), and such a photoionization mechanism would not be expected to show a
significant KIE if the coupling mode involved in electron transfer does not involve O–H(D) stretching.\textsuperscript{43,64}

The autoionization mechanism that we have proposed for photoexcited phenol has some similarities with phenolate (PhO\textsuperscript{−}) in water, which generates phenoxy radical and solvated electrons within the first few picoseconds.\textsuperscript{22} The far more rapid autoionization observed for PhO\textsuperscript{−}, likely reflects the 0.7 eV lower vertical ionization potential.\textsuperscript{30} The phenol and phenolate autoionization reactions yield the same products, however, ionization of phenol also produces a third product, a proton. This complicates the geminate recombination pathways, and allows for many possible sequential reactions (see SI). The two reactions that will most probably affect the yield of phenoxy radical and solvated electron transients in our data, include the same recombination reaction as for phenolate: phenoxy radicals and electrons combine to form ground state phenolate molecules, that subsequently abstract a proton from an adjacent water molecule. The other mechanism involves a proton scavenging a solvated electron and then upon encounter with the phenoxy radical, reformation of the O–H bond.

The present time-resolved TA data for phenol in aqueous H\textsubscript{2}O and D\textsubscript{2}O solutions are therefore most plausibly understood in terms of photoinduced autoionization. Identical photoproducts are observed at short (200 nm) and long (267 nm) excitation wavelengths, but with very different formation timescales (femtosecond \textit{vs.} nanosecond, respectively). Electrons are ejected immediately following 200 nm excitation; the high-energy phenol radical cation intermediate is formed and rapidly deprotonates to yield PhO\textsuperscript{−} radicals. In contrast, at 267 nm the photoexcited phenol molecules initially relax to the PhOH(S\textsubscript{1(v=0)}) level, \textit{i.e.} to energies just above the minimum required to produce a PhO\textsuperscript{−} radical, a solvated electron and a proton (see energetics in the SI). The measured timescales for forming these products allied with the lack of measurable
H/D KIE encourages the view that the observed products are formed via near-threshold autoionization, though the involvement of an excited state PCET mechanism cannot be wholly excluded at this stage.

Phenol in aqueous solution surely constitutes one of the simpler test beds for \textit{ab initio} calculations designed to explore near-threshold autoionization or possible excited state PCET mechanisms.\textsuperscript{65,66} Understanding the nuclear motions associated with the solute, and any coupling to specific bath modes, are now crucial in order to reveal any role for PCET at low excitation energies and to advance a detailed understanding of our favored mechanism – autoionization – in phenol, and in larger biomolecules in which phenol (and related sub-units) acts as a chromophore.

ASSOCIATED CONTENT

\textbf{Supporting Information} Experimental summary, transient absorption spectra at 200 nm in the presence and absence of H\textsuperscript{+} electron scavengers, schematic free energy landscape associated with gas phase and solvated phenol, transient absorption spectra for 267 nm irradiation of phenol in with and without H\textsuperscript{+} measured at $t < 1$ ns. This material is available free of charge via the Internet at \url{http://pubs.acs.org}.

AUTHOR INFORMATION

\textbf{Corresponding Author}

*Email: stephen.bradforth@usc.edu
ACKNOWLEDGMENTS

The authors would like to thank Gaurav Kumar, Saptaparna Das at the University of Southern California, Stephanie Harris and Daniel Murdock at the University of Bristol for their insights and useful discussions. Research at USC was supported by the US National Science Foundation under CHE-1301465 and at Bristol by the Engineering and Physical Sciences Research Council via Programme Grant EP/G00224X.

REFERENCES


(43) Pino, G. A.; Oldani, A. N.; Marceca, E.; Fiji, M.; Ishiuchi, S. I.; Miyazaki, M.;


