Reconstructing the structure of directed and weighted networks of nonlinear oscillators

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The formalism of complex networks is extensively employed to describe the dynamics of interacting agents in several applications. The features of the connections among the nodes in a network are not always provided beforehand, hence the problem of appropriately inferring them often arises. Here, we present a method to reconstruct directed and weighted topologies (REDRAW) of networks of heterogeneous nonlinear oscillators. We illustrate the theory on a set of representative examples.

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I. INTRODUCTION

The study of complex networks provides a rigorous formalism to describe phenomena involving populations of interacting agents in fields as diverse as Physics, Engineering, Biology, Chemistry, Social Science and the Internet.

Often in applications, the topology of the interactions (links) among the agents (nodes) in a network is not known a priori. In these cases the structure of the network must be reconstructed from measurements of the nodes’ dynamics; a problem often known in the literature as network reconstruction or network inference. The need for network inference strategies arises in several contexts, e.g., reconstructing functional activation/repression links in gene regulatory networks or causal relationships in stochastic processes, understanding the structure of social interactions in a group from communication data, and inferring functional relationships between areas in the brain from EEG data or in physiological systems.

A notable case of interest is that of reconstructing the structure of networks whose nodes exhibit oscillatory dynamics (e.g., neurons, cellular cycles, synthetic biological oscillators, groups of walking autonomous robots). For these networks, several reconstruction approaches have been proposed. Examples include the methodology presented in [18], dealing with the problem of inferring directed and weighted topologies of networks of Hindmarsh-Rose neurons and Lorenz oscillators; the strategy discussed in [19], investigating community detection in undirected and unweighted networks of Kuramoto oscillators; the techniques developed in [21], addressing the issue of reconstructing directed topologies of small networks of coupled phase oscillators; and the methodology presented in [22], dealing with the reconstruction of directed weighted networks of coupled phase oscillators driven by temporarily constant input signals. Recent approaches to inferring the structural connectivity of a network from its dynamics have also been reviewed in [23].

Previous methods suffer from several drawbacks. Often they can only reconstruct undirected networks and do not include an ex post analysis of the reconstructed topology to isolate the presence of false positives (reconstructed links that do not exist in the real network), or if they do, no systematic method is presented on how to select cut-off thresholds to remove them [19]. The only exceptions are the works presented in [11, 24]. However, such methodologies require the observed dependency matrix to be either invertible or diagonalizable, and have order $O(n^3)$ of computational complexity (with $n$ being the number of nodes in the network), as opposed to our algorithm which does not make any assumption on the structure of the inferred interactions and has lower computational complexity (as discussed later in the text).

The aim of this work is to present REDRAW (REconstruciton of DiRected And Weighted topologies), a novel method for reconstructing directed and weighted topologies of networks of nonlinear oscillators. REDRAW radically extends the strategy in [19], which was originally conceived to only detect communities in undirected and unweighted networks of Kuramoto oscillators [25, 26]. In what follows, we first describe the approach including a systematic algorithm to perform the ex post analysis of the reconstructed network and select appropriate cut-off thresholds to remove false positives. We show that REDRAW allows to effectively infer directed and weighted links among an ensemble of coupled heterogeneous oscillators. We then validate the method on a number of representative examples, including a set of real-world networks obtained from [27-31], highlighting its advantages and limitations.

II. MATERIALS AND METHODS

We start by considering a network of $n$ nonlinearly coupled heterogeneous oscillators described by:
\[
\dot{x}_i = f_i(x_i) + \frac{c}{n} \sum_{j=1}^{n} a_{ij} g(x_i, x_j), \; i = 1, 2, \ldots, n \quad (1)
\]

where \(x_i \in \mathbb{R}^p\) is the \(p\)-dimensional state of the \(i\)th oscillator, \(f_i\) denotes its dynamics, \(g\) is a generic nonlinear coupling function, \(c > 0\) represents the global coupling strength among all nodes in the network, and \(a_{ij}\) represents the local influence that node \(j\) has on node \(i\). For directed and weighted topologies, in general \(a_{ij} \neq a_{ji}\) with \(a_{ij} \geq 0 \forall i, j\). By means of standard techniques \[20\], the system in Eq. (1) can be reduced to a network of nonlinear phase-coupled oscillators:

\[
\dot{\theta}_i = \omega_i + \frac{c}{n} \sum_{j=1}^{n} a_{ij} \Gamma(\theta_j - \theta_i), \; i = 1, 2, \ldots, n \quad (2)
\]

where \(\theta_i \in [-\pi, \pi]\) represents the phase of the \(i\)th oscillator, \(\omega_i > 0\) its natural frequency, and \(\Gamma(\theta_j - \theta_i)\) is a generic 2\(\pi\)-periodic function.

We assume that a data-set of \(K \times n\) time series of length \(T\) is available. This is obtained by acquiring \(n\) time series of duration \(T > 0\) (one for each of the \(n\) oscillators in the network of interest) during \(K\) experiments (or simulations) where nodes are started from different initial conditions and phase-locking is achieved \[32\].

The goal of REDRAW is to reconstruct the topology of the network by inferring, from the data, the presence and directions of links between the oscillators. The method also captures the relative strength of the interactions among nodes by establishing what links are stronger than others. However, it does not aim at estimating the precise value of the edge weights. Note that no perfect knowledge of the systems dynamics is necessary for our method. The only assumption is that the measured phase differences among the oscillators give an indication of the strength and direction of their mutual influence.

REDRAW is based on the following six steps.

**Step 1.** The relative phase \(\Delta \theta_{ij,k}(t) = \theta_{i,k}(t) - \theta_{j,k}(t) \in [-\pi, \pi], \forall k = 1, \ldots, K\) with \(t \in [0, T]\) is evaluated for every pair of nodes in each of the \(K\) experiments. Negative values of \(\Delta \theta_{ij,k}(t)\) indicate that node \(i\) lags behind node \(j\) at time instant \(t\) of the \(k\)th experiment.

**Step 2.** The parameter \(\zeta_{ij,k}(t)\), representing the influence of node \(j\) on node \(i\) at time \(t\) of the \(k\)-th experiment, is defined as:

\[
\zeta_{ij,k}(t) := \begin{cases} 
1 + \frac{h[\Delta \theta_{ij,k}(t)]}{2}, & \Delta \theta_{ij,k}(t) \leq 0 \\
0, & \Delta \theta_{ij,k}(t) > 0
\end{cases} \quad (3)
\]

and is calculated from the available data for all pairs of nodes in the network, where \(h\) is an even function in the interval \([-\pi, \pi]\) characterized by the following properties: \(ij\) it takes its maximum value 1 in the origin [e.g., \(h(0) = 1\)]; \(ii\) it takes its minimum values \(-1\) at the extremes of its domain [e.g., \(h(-\pi) = h(\pi) = -1\)]; \(iii\) it is monotonically increasing in the left-half interval \([-\pi, 0]\) and monotonically decreasing in the right-half interval \([0, \pi]\). Note that \(\zeta_{ij,k}(t) \neq \zeta_{ji,k}(t)\) and that \(\zeta_{ij,k}(t) \in [0, 1]\), where \(\zeta_{ij,k}(t) = 1\) represents the maximum level of influence that node \(j\) has on node \(i\). When \(\Delta \theta_{ij,k}(t) > 0\), agent \(j\) lags behind agent \(i\), hence the influence that the former has on the latter is set to 0.

**Step 3.** The time-average of the influence parameter defined in Eq. (3) is evaluated for each pair of nodes as:

\[
\rho_{ij,k} := \frac{1}{T} \int_{0}^{T} \zeta_{ij,k}(t) \, dt \quad (4)
\]

**Step 4.** The weight of the link directed from node \(j\) to node \(i\) is computed for every pair of nodes by averaging the quantity defined in Eq. (4) over the total number of experiments as:

\[
\rho_{ij} := \frac{1}{K} \sum_{k=1}^{K} \rho_{ij,k} \quad (5)
\]

**Step 5.** Links whose \(\rho_{ij}\) is null are discarded from the reconstructed network. Non-zero weights \(\rho_{ij}\) are checked among all triplets of connected nodes so that, on the basis of their intensity, one can be possibly regarded as a false connection and set to zero (the corresponding link is then removed). As an example, consider a triplet of three connected nodes, say \(w, y\) and \(z\) (Fig. 1). If the weight \(\rho_{zw}\) between \(w\) and \(z\) is lower than a certain threshold \(0 \leq \nu < 1\) and is also lower than both the weight \(\rho_{yw}\) between \(w\) and \(y\), and the weight \(\rho_{yz}\) between \(y\) and \(z\), then \(\rho_{zw}\) is set to 0 and the link between the pair \((w, z)\) is removed as it can be regarded as indirect effect of \(w\) on \(z\) through node \(y\). This is a modified version of the standard Data Processing Inequality (DPI) often used in network reconstruction methods \[33\]. Note that the higher the value of \(\nu\), the less connected triplets are found in the network.

**Step 6.** Network Thresholding: all parameters \(\rho_{ij}\) whose value is below a certain threshold \(\mu\) are set to 0.
with $0 \leq \mu \leq \nu$ (the higher the value of $\mu$, the sparser is the reconstructed network structure).

Four standard metrics are used to assess the performance of REDRAW [34]. They are the Positive Predicted Value (PPV), the Accuracy Rate (ACC), the True Positive Rate (TPR) and the False Positive Rate (FPR). Denoting with $N_{TP}$, $N_{FP}$, $N_{TN}$, $N_{FN}$ and $N_{TOT} := n(n - 1)$ the total number of true positives, false positives, true negatives, false negatives and possible links among all the nodes in the network, respectively, such metrics are defined as:

$$PPV := \frac{N_{TP}}{N_{TP} + N_{FP}}, \quad ACC := \frac{N_{TP} + N_{TN}}{N_{TOT}},$$

$$TPR := \frac{N_{TP}}{N_{TP} + N_{FN}}, \quad FPR := \frac{N_{FP}}{N_{FP} + N_{TN}}.$$

Since the interactions between $n(n - 1)$ pairs of nodes have to be inferred for each available experiment, REDRAW has order $O(Kn^2)$ of computational complexity. On the other hand, since the performance metrics are independent of the number of numerical experiments available, their application has order $O(n^2)$.

**Remark 1** If phase-locking is not achieved by the oscillators in the network, either no links or few weak links are inferred according to Eq. (6). Therefore, the influence that the phase of each node has on that of the others is negligible, hence the strength of the interactions among the oscillators is not sufficiently high for them to form any clusters.

**Remark 2** The time interval $[0,T]$ can be partitioned in $L$ time windows $[t_l,t_{l+1}]$ of length $\Delta T_l$, with $l = 0, \ldots, L - 1$ so that the network topology can be reconstructed over each subinterval to show how the nodes’ interactions evolve over time. The choice of $\Delta T_l$ is dependent on the specific application of interest (e.g. nodes’ temporal dynamics).

### III. RESULTS

To validate and illustrate REDRAW we consider a network of $n$ nonuniform Kuramoto oscillators [35, 36], employed as a model generating the data to test the algorithm. It is obtained by setting $\Gamma(\theta_j - \theta_i) = \sin(\theta_j - \theta_i - \phi_{ij})$ as an instance of Eq. (2), described by:

$$\dot{\theta}_i = \omega_i + \frac{c}{n} \sum_{j=1}^{n} a_{ij} \sin(\theta_j - \theta_i - \phi_{ij}), \quad i = 1, 2, \ldots, n$$

(6)

The phase shift

$$\phi_{ij} := \begin{cases} \frac{\phi_{ij}}{a_{ij}} , & a_{ij} > 0 \\ 0 , & a_{ij} = 0 \end{cases}$$

(7)

represents how much node $i$ lags behind node $j$, with $\phi_{ij} \in \left[0, \frac{\pi}{2}\right]$. Note that the higher the influence $a_{ij}$ that node $j$ has on node $i$, the lower the phase shift $\phi_{ij}$ (i.e., the less node $i$ lags behind node $j$). Furthermore, we choose $h(\theta) = \cos(\theta)$ in Eq. (3).

In contrast to previous methods, we provide guidelines on the selection of the thresholds $\nu$ and $\mu$ in Steps 5 and 6, by performing a systematic analysis of how the performance metrics change as the thresholds are varied when inferring *in-silico* known ensembles (in general different from the original system to be reconstructed) of nonlinear oscillator networks [37]. Specifically, we applied the algorithm to reconstruct $N$ randomly generated graphs following the Erdős-Rényi model [38, 39] as well as networks generated using the Barabási-Albert method [40], for ensembles of nonuniform Kuramoto oscillators of different size $n$. Note that any topology and oscillatory dynamics could be employed, however for the latter it is sensible to choose the same model as that of the system it is desired to reconstruct when information is given on the dynamics of its nodes. For both models, we evaluated what average values of $\nu$ and $\mu$ lead to acceptable reconstruction metrics, and propose that these values should then be selected when inferring an unknown topology (Fig. 2).

**Remark 3** If the uncoupled dynamics is much larger than the coupling [e.g., $\omega_i \gg \frac{c}{n}$ in Eq. (6)], some oscillators may not be phase-locked with their neighbors even though they share a connection. Therefore, in the following numerical examples we assume that the natural frequencies of the oscillators are all bounded.

We start by applying our method to reconstruct the four different topologies in Figs. 3(a)-3(d). Specifically, the values of the natural frequencies $\omega_i$ were randomly drawn from the standard uniform distribution on the interval $[1, 2]$ rad s$^{-1}$, and so were the initial conditions $\theta_i(0)$ from the interval $[-\pi, \pi]$, $i = 1, \ldots, 4$. The model parameters were set to $c = 10$ and $\phi = \frac{\pi}{2}$ so that phase-locking could be achieved. For each topology, a synthetic dataset of $K = 50$ experiments of duration $T = 30s$ was obtained assuming the structure of the network to be unknown. Thresholds were set to $\nu = 0.9$ and $\mu = 0.7$ from Figs. 2(a-c).

The reconstructed networks are represented in Figs. 3(e)-3(h), respectively. For each topology, neither missing links nor false positives are found, and the directionality of the links is correctly inferred for all of them ($PPV = ACC = TPR = 100\%$, $FPR = 0\%$). As for the weights, their magnitudes’ relative relationship in the assigned topologies is correctly inferred as well. For instance, note how $a_{34} < a_{23} < a_{12}$ in Fig. 3(b) corresponds to $\rho_{34} < \rho_{23} < \rho_{12}$ in Fig. 3(f), or how $a_{31} = a_{23} < a_{12}$ in Fig. 3(c) corresponds to $\rho_{31} = \rho_{41} < \rho_{21}$ in Fig. 3(g).

The evolution over time of the reconstructed network with topology represented in Fig. 3(c) was inferred over time windows of length $\Delta T_l = \Delta T = 0.5s$, and is shown...
FIG. 2. Selection of thresholds for networks of $n$ nonuniform Kuramoto oscillators. Acceptable ranges within which $\nu$ (x–axis) and $\mu$ (y–axis) should take values when reconstructing an unknown network of $n = 4$ (a,c) or $n = 20$ (b,d) oscillators are depicted as shades-of-grey maps; $m$ denotes the number of metrics for which acceptable values are achieved when in-silico reconstructing known test topologies. Specifically, (a) and (b) refer to the metrics obtained when employing random graphs following the Erdős–Rényi model as test topologies, whereas (c) and (d) refer to those obtained for networks generated using the Barabási–Albert method. The parameters of the oscillators were set so as to achieve phase-locking. The lighter the color of the area where $(\nu, \mu)$ belongs, the higher $m$, hence the more sensible the choice of the thresholds. Ideally, the thresholds to be used when inferring unknown topologies should be selected within white regions where $\nu$ and $\mu$ lead to acceptable values for all the four metrics here employed. See [37] for more details.

FIG. 3. Assigned and inferred topologies, $n = 4$. The topologies (a)-(d) represent those used in the numerical simulations to generate the data-set then employed to test our method. The reconstructed topologies are respectively depicted in the bottom panels (e)-(h). The numerical values on the edges of the assigned topologies represent the values of $a_{ij}$ in the model described in Eq. (6), whereas the italic numerical values on the edges of the inferred topologies represent parameters $\rho_{ij}$ estimated by REDRAW.

Analogous results are found for the other topologies (data not shown).

Next, we tested REDRAW on larger, more challenging networks of $n = 17$ nodes. The structures being considered were obtained by interconnecting four sub-networks with the structure represented in Fig. 3(d), either through a central hub in a geometric graph configuration [32], or as the Ravasz-Barabási graph [3]. Details on the results obtained for the reconstruction of such topologies can be found in [41].

As a further validation of the methodology, we used REDRAW to reconstruct the networks of $n = 20$ nodes shown in Fig. 4 both assumed to be unknown. The coupling strength among nodes was set as $c = 50$, and the thresholds of the algorithm to $\nu = 0.65$ and $\mu = 0.60$ from Figs. 2(b,d). All the other parameters were selected as before.

We find that the directionality and weights of all edges in the regular network (Fig. 4(a)) are correctly inferred [Fig. 4(c)], with the only exception of a missing link from node 2 to node 4.

The introduction of 5 long-distance edges [Fig. 4(b)] is well captured by the topology inferred in Fig. 4(d), leading to the formation of an equal number of clusters [nodes (1–4), (5–8), (9–12), (13–16) and (17–20)]. For each of the 5 nodes being influenced by one of the additional links (nodes 4, 8, 12, 16, 20), no outgoing links are inferred. This is a result of the model described in Eq. 6: for instance, the phase of node 4 is lowered by the influence of node 19, thus leading to higher mismatches of the former with nodes 5 and 6 (similar reasoning can be carried out for nodes 8, 12, 16 and 20). However, note that such edges can be inferred when selecting lower values for the thresholds.

Lastly, we applied REDRAW to reconstruct some real-world network topologies from the literature [27–31], and explored the effects of noise and uncertainty on the measured phases [43].

The values of the four reconstruction metrics are detailed in Table I for the various topologies considered in this work, together with their structural parameters, the value of the coupling strength $c$ used in Eq. 6 and selected in order for the network to achieve phase-locking, and that of the thresholds $\nu$ and $\mu$ used for the reconstruction and selected according to [37]. All other parameters were kept as in the previous examples. Assuming $PPV$, $ACC$, $TPR > 50\%$ for the first three metrics and $FPR < 10\%$ for the last one as acceptability criteria, Table I shows the effectiveness of REDRAW in inferring networks of different size (the criteria are met in 91% of the cases).

For all the topologies considered in this work, we quantified the performance metrics for different numbers of experiments and observed how, for a sufficiently high number of repetitions, REDRAW is not sensitive to the specific value of $K$. For the sake of brevity, we reported here only results for the case $K = 50$; in fact, when $K \geq 50$, possible fluctuations of the metrics are negligible. For more details on the performance of REDRAW for different number of experiments see [44].

Remark 4 Suppose an external field drives a set of phys-
FIG. 4. Assigned and inferred topologies, $n = 20$. The regular network (a) and that obtained by some long-distance rewiring (b) on the left-hand side represent the topologies used in the numerical simulations to generate the data-set then employed to test our method. The reconstructed topologies are respectively depicted on the right-hand side (c,d). Different scales of gray quantify the numerical value of $a_{ij}$ for the assigned topologies, and $\rho_{ij}$ estimated by REDRAW for the inferred ones.

TABLE I. Validation of REDRAW. This table shows the four performance metrics obtained for different topologies, together with their own number of nodes $n$ and edges $e$, respectively. The values of coupling strength $c$ and thresholds $\nu$ and $\mu$ employed in the numerical simulations are detailed as well.

<table>
<thead>
<tr>
<th>Topologies</th>
<th>n</th>
<th>e</th>
<th>c</th>
<th>$\nu$</th>
<th>$\mu$</th>
<th>PPV(%)</th>
<th>ACC(%)</th>
<th>TPR(%)</th>
<th>FPR(%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Enzyme-Catalyzed reaction pathway [27]</td>
<td>8</td>
<td>7</td>
<td>20</td>
<td>0.90</td>
<td>0.75</td>
<td>100</td>
<td>100</td>
<td>100</td>
<td>0</td>
</tr>
<tr>
<td>Songbird brain [28]</td>
<td>12</td>
<td>13</td>
<td>35</td>
<td>0.90</td>
<td>0.65</td>
<td>63</td>
<td>93</td>
<td>77</td>
<td>5</td>
</tr>
<tr>
<td>Bank stocks connections [29]</td>
<td>16</td>
<td>32</td>
<td>40</td>
<td>0.80</td>
<td>0.70</td>
<td>56</td>
<td>87</td>
<td>67</td>
<td>4</td>
</tr>
<tr>
<td>Regular network</td>
<td>20</td>
<td>40</td>
<td>50</td>
<td>0.65</td>
<td>0.60</td>
<td>100</td>
<td>99</td>
<td>97</td>
<td>0</td>
</tr>
<tr>
<td>Rewired network</td>
<td>20</td>
<td>45</td>
<td>60</td>
<td>0.65</td>
<td>0.60</td>
<td>67</td>
<td>92</td>
<td>67</td>
<td>4</td>
</tr>
<tr>
<td>Human PPI [30]</td>
<td>23</td>
<td>22</td>
<td>40</td>
<td>0.80</td>
<td>0.50</td>
<td>100</td>
<td>98</td>
<td>45</td>
<td>0</td>
</tr>
<tr>
<td>Human PPI [30]</td>
<td>25</td>
<td>27</td>
<td>50</td>
<td>0.80</td>
<td>0.65</td>
<td>56</td>
<td>96</td>
<td>44</td>
<td>2</td>
</tr>
<tr>
<td>Hainan Power Grid Company [31]</td>
<td>48</td>
<td>63</td>
<td>150</td>
<td>0.60</td>
<td>0.60</td>
<td>52</td>
<td>97</td>
<td>62</td>
<td>2</td>
</tr>
</tbody>
</table>

DPI techniques of Step 5.

IV. CONCLUSION AND DISCUSSION

We presented an effective reconstruction method to infer directed and weighted topologies of networks of heterogeneous nonlinear oscillators from data on their dy-
namics. We tested the strategy on a number of different examples, showing that in all cases the methodology guarantees good accuracy and an extremely low number of false positives, which might occur when physically disconnected nodes are coincidentally characterized by similar oscillation frequencies and phases.

REDRAW can be applied to reconstruct the interactions among any nodes exhibiting oscillatory dynamics, as long as their phase trajectories are available. In this work, we tested the algorithm on networks of Kuramoto oscillators only as a case-study. Other possibilities include, for instance, networks of Lorentz or Rössler oscillators, whose phases can be estimated by means of the techniques proposed in [45, 46].

Currently, we are applying the algorithm to reconstruct networks of social interactions in groups of individuals moving in synchrony [47–49]. Such results are beyond the scope of this work and will be presented elsewhere.

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[44] See section 3 of Supplementary Material at LINK for detailed values of the standard metrics as a function of the number of experiments $K$ (2016).